

Appendix: Elements of probability theory.

(A1)

see the book of Chris hope and also math lecture

→ We have seen the necessity of principles to use a stat. description

◦ practical necessity.

◦ Absence of conserved q'ty → dynamics becomes chaotic, i.e. dependent on exact initial conditions that we do not know!

◦ Actually, we do not know the exact Hamiltonian either → isolated system do not really exist.

◦ something I did not mention, is that we ~~can~~ ^{shall} also use Q. Mechanic to describe our system → inherent probabilist description

I) Random variables

1) Discrete one

→ We consider a variable X taking the value $x^{(m)}$ if the event e_m is realized. X is a random variable. We can define the probabilities $\{P_m\}$ such that $P(x) \equiv \text{Proba}(X=x^m) = \sum_{m=1}^M P_m \delta_{x, x^{(m)}}$

→ Here we assumed the space of events $\mathcal{E} = \{e_1, \dots, e_M\}$ M possible values

Toss a coin → head or queue $M=2$

consider N realization of the experiment. We obtain N_m the event e_m

$N_m \leq N$ is random

$$P_m \stackrel{\text{def}}{=} \lim_{N \rightarrow \infty} \frac{N_m}{N} \Rightarrow \sum_{m=1}^M P_m = 1$$

Ex: Jackpot -

x	β	γ
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with $x, \beta, \gamma \in [0, 9]$

1000 events $P_{x\beta\gamma} = \frac{1}{1000}$

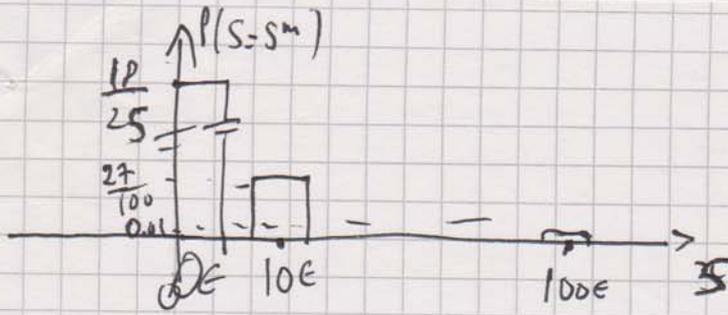
S = sum of gain → random variable → 3 values

Assume if $x = \beta = \gamma \rightarrow$ win $100€ = S_1$
else $0 = S_3$
- $x \neq \beta$ or $x \neq \gamma$ or $\beta \neq \gamma \rightarrow 10€ = S_2$

$$P(S_1) = \frac{10}{1000} \text{ for } S_1 = 100€$$

$$P(S_2) = \frac{270}{1000} \text{ for } S_2 = 10€$$

$$P(S_3) = \frac{720}{1000} = \frac{18}{25} \text{ for } S_3 = 0€$$



2) Continuous variable

In the previous example $S \in \{S^1, S^2, S^3\} \rightarrow$ discrete set

In some case the random variable can be \subseteq

ex: brownian motion of a microscopic object at the surface of a liquid \rightarrow Jean Perrin, 1908

$\vec{r} = (x, y)$ where both x, y are random positi^o

The extension to a \subseteq variable is as follows:

we introduce the probability density $w(x)dx$ that the random variable x takes the value x within dx :

$$w(x) dx \equiv \text{Proba} \{x \in [x, x+dx]\}$$

$$\int w(x) dx = 1 \quad \text{normalisat}^o$$

change of variable: If Y random variable such that $X = f(Y)$
 $P(y) dy = ?$ f bijective

$$P(y) dy = \text{Proba} \{X \in [f(y), f(y) + f'(y) dy]\}$$

$$= w(f(y)) f'(y) dy$$

Thus $P(y) = w(f(y)) f'(y)$

Let us check the normalisat^o $\int P(y) dy = \int w(f(y)) f'(y) = \int w(x) dx = 1$

Ex $E = \frac{p^2}{2m}$
 $\frac{d}{dx}$

suppose we know the distribut^o of p say $P(p)$

then $P(E) dE = P(p) dp$
 $P(E) \frac{d}{m} dp = P(p) dp$

$$\mathcal{P}(p) dp = w(E) dE = \mathcal{P}(\sqrt{2mE}) \sqrt{\frac{m}{2E}} dE = w(E) dE$$

(A3)

$$w(E) = \mathcal{P}(\sqrt{2mE}) \sqrt{\frac{m}{2E}}$$

If we have N random variables $w(x_1, \dots, x_N) = \text{Prob}\{x_1 \in [x_1, x_1 + dx_1], \dots, x_N \in [x_N, x_N + dx_N]\}$

II) Average, variance, moments, cumulants & generating function

We consider the probabilistic average. (this assumes that the number of realizations of an experiment, say N , is large $\rightarrow N \rightarrow \infty$)

noted $E(X)$ in math

$$\langle X \rangle = \sum_{n=1}^M P_n X^{(n)} \quad \text{for discrete variables}$$

In our roulette game: $M = 3000$, However S can take only 3 values

$$\langle X \rangle = \sum_{v=1}^V P_v x_v \quad \leftarrow \text{here the sum.}$$

$$\langle S \rangle = 100 \times \frac{1}{100} + 10 \times \frac{27}{100} + 0 \times \frac{18}{15} = 3.70 \text{€}$$

The casino has to make sure that the cost of the try is > 3.70 to avoid bankruptcy.

Continuous variable $\langle X \rangle = \int dx w(x) x$

and $\langle f(X) \rangle = \int dx w(x) f(x)$

Variance

$$\text{var}(X) = \langle (X - \langle X \rangle)^2 \rangle = \langle X^2 \rangle - \langle X \rangle^2$$

standard deviation $\sigma_x = \sqrt{\text{var}(X)}$

characterizes the typical standard deviation from a random variable around its mean value

If $w(x) = \frac{1}{a} e^{-x/a}$ for $x > 0$

$$\langle x \rangle = \int_0^{\infty} dx \frac{x}{a} e^{-x/a} = a \int_0^{\infty} dy y e^{-y}$$

$$\langle x^2 \rangle = 2a^2 = a \left[-y e^{-y} \right]_0^{\infty} + a \int_0^{\infty} dy e^{-y} = a$$

(A4)

Add

~~$\Delta x = a =$~~

We can extend these definitions to higher moments

Correlat^o (SIDE Page)

$$\mu_n \equiv \langle x^n \rangle = \int dx x^n w(x)$$

RQ: If we know all $\mu_n \forall n$, then $w(x)$ is known provided $w(x)$ is non pathological.

If $\sigma_x \ll |\langle x \rangle|$ then $\langle x^2 \rangle = \langle x \rangle^2 + \text{var} X \approx \langle x \rangle^2$

Basically the second moment is given by the square of μ_1 => the variance brings some extra information compared to μ_1

We also call the variance $\left. \begin{matrix} \mu_1 = c_1 \\ \sigma_x^2 \equiv c_2 \end{matrix} \right\} \rightarrow$ second cumulant

We can also extend the notion of cumulants to larger n with the idea that c_n brings new info not contained in $\{c_i \in n\}$

1) $c_1 \equiv \mu_1 = \langle x \rangle$

2) $c_2 \equiv \mu_2 - \mu_1^2 = \sigma_x^2 = \langle \tilde{x}^2 \rangle$ where $\tilde{x} = x - \langle x \rangle$

=> width of the distribut^o

→ picture

3) skewness = $\langle (x - \langle x \rangle)^3 \rangle$

$$c_3 = \langle \tilde{x}^3 \rangle = \mu_3 - 3\mu_2\mu_1 + 2\mu_1^3$$

=> asymmetry of the distribut^o function

$c_3 > 0 \rightarrow$ fluctuations are positive

4) Kurtosis ~~is~~ $\langle \tilde{x}^4 \rangle - 3 \langle \tilde{x}^2 \rangle^2$

For a Gaussian distrib^o funct^o $\langle \tilde{x}^4 \rangle = 3 \langle \tilde{x}^2 \rangle^2$

$$c_4 = \langle \tilde{x}^4 \rangle - 3 \langle \tilde{x}^2 \rangle^2 \Rightarrow \text{deriv^o from Gaussian}$$

show plots

$c_4 > 0 \Rightarrow$ large fluctu^o are more favorable

Q: \rightarrow Is there a systematic way to expand cumulants?

\Rightarrow Generating functions

$$g(p) \equiv \langle e^{-px} \rangle = \int dx e^{-px} w(x) = \sum_n \frac{(-p)^n}{n!} \mu_n$$

Laplace transform of the proba distrib^o

$$g(p) = \sum_{n=0}^{\infty} \frac{(-p)^n}{n!} \mu_n \Rightarrow \mu_n = (-1)^n g^{(n)}(p) \Big|_{p=0}$$

$$\mu_n = (-1)^n \frac{d^n g(p)}{dp^n} \Big|_{p=0}$$

\Rightarrow Easier to calculate $\langle e^{-px} \rangle$ than all μ_n !

Let us define the function $w(p) \equiv \ln g(p)$

$$w(p) = \sum_{n=1}^{\infty} \frac{(-p)^n}{n!} c_n \Rightarrow c_n = (-1)^n \frac{d^n w(p)}{dp^n} \Big|_{p=0}$$

$$\ln g(p) \approx \ln(1 - \mu_1 p + \frac{1}{2} \mu_2 p^2 + O(p^3))$$

$$= -\mu_1 p + \frac{1}{2} (\mu_2 - \mu_1^2) p^2 + O(p^3) \Rightarrow \begin{cases} c_1 = \mu_1 \\ c_2 = \mu_2 - \mu_1^2 \\ \dots \end{cases}$$

III Some standard statistical laws

ex: binomial laws

N events $\begin{cases} \rightarrow$ tail with proba p \Rightarrow what is the proba to get n tails \\ \rightarrow face \rightarrow 1-p \rightarrow head

$$(s = e^{-p}) \quad \Pi_N(n) = C_N^n p^n (1-p)^{N-n}$$

$$g_N(s) \equiv \langle s^n \rangle \Rightarrow \text{we find } g_N(s) = [1 + p(s-1)]^N$$

indeed $g_N(s) = \sum_n \Pi_N(n) s^n = \sum_n C_N^n (ps)^n (1-p)^{N-n} = (1-p + ps)^N$

Poissonian law \Rightarrow Applies for independent events

Proba of an event to occur during time $dt = \lambda dt$

\Rightarrow Applies to Schottky exp \rightarrow proba of extracting an el. under a bias
 $\langle I \rangle = \frac{1}{T} \int_0^T \lambda e dt = \lambda e$

$P_{\lambda T}$ of having n events in a time T

$P_{\lambda T = n} = \frac{\lambda^n}{n!} e^{-\lambda T}$ with $\lambda T = \mu$ Poisson law

$g_{\mu}(p) = \exp(\mu(e^p - 1))$ and $w_{\mu}(p) = \mu(e^p - 1)$

$c_n = (-1)^n \frac{d^n}{dp^n} w_{\mu}$ $= \mu \sum_{n=1}^{\infty} \frac{(-p)^{n-1}}{(n-1)!}$

$\Rightarrow c_n = \mu \quad \forall n$

\Rightarrow all the cumulants are equal!

Ex: Emission of e^- in a tunnel junction \Rightarrow $\left\{ \begin{array}{l} \neq$ bunching (photons) \\ \neq antibunching corr $e^- \rightarrow$ Pauli
• Radiation decay \rightarrow SLIDES

Gaussian law

$P_{a,\sigma}(x) = \frac{1}{\sqrt{2\pi}\sigma} e^{-\frac{(x-a)^2}{2\sigma^2}}$ \rightarrow 2 parameters $\left\{ \begin{array}{l} a \\ \sigma \end{array} \right.$

$g_{a,\sigma}(p) = \int dx P_{a,\sigma}(x) e^{ipx}$
 $= \exp(-ap + \frac{1}{2} \sigma^2 p^2)$

$\Rightarrow w_{a,\sigma}(p) = -ap + \frac{1}{2} \sigma^2 p^2 \rightarrow \left\{ \begin{array}{l} c_1 = a = \langle x \rangle \\ c_2 = \text{Var}(x) = \sigma^2 \end{array} \right.$

$c_{n>2} = 0$

\Rightarrow There might exist pathological distributions

$w(x) \sim |x|^{-\mu-1}$ when $|x| \rightarrow \infty \rightarrow \mu_n$ finite iff $n < \mu$

Ex: Cauchy law $w(x) = \frac{a/\pi}{a^2 + x^2}$

AA

here $\mu=1 \Rightarrow$ all moments are ∞

$\tilde{g}(k) = \langle e^{-i\frac{k}{2}x} \rangle = e^{-a|k|} \Rightarrow$ no Taylor expansion in $h=0$

IV) Central limit theorem

We consider a set of N random variables X_1, \dots, X_N

- Assume:
- i) They are statistically independent
 - ii) They all follow the same distribut^o function $p(x)$
 - iii) We only assume $\langle X_n \rangle$ and $\langle X_n^2 \rangle$ are finite i.e. the two first moments do exist!

Question: What is the distribution of $S = \sum_{i=1}^N X_i$?

RQ: • Up to a prefactor $\frac{1}{N}$ S coincides with the statistical average
 $F^{(N)} = \frac{1}{N} \sum_{i=1}^N X_i \Rightarrow S$ therefore controls the statistical error made by n^{th} sample \Rightarrow vote

• GAZ $E_c \gg E_p \Rightarrow E \sim E_c = \sum_{i=1}^N \frac{p_i^2}{2m} \Rightarrow$ Assuming p_i , R.V E_c ?

Look at the generating function:

$G_N(p) \equiv \langle e^{-pS} \rangle = \langle \exp(-p \sum_{n=1}^N X_n) \rangle$

$g(p) = \langle e^{-px} \rangle$
 $w(p) = \ln g(p)$

$\stackrel{\text{hyp 1}}{\Rightarrow} = \langle \prod_{n=1}^N (e^{-pX_n}) \rangle$ statistical independent
 $= \prod_{n=1}^N \langle e^{-pX_n} \rangle = |g(p)|^N$
 $\stackrel{\text{hyp 2}}{\Rightarrow}$

$\Rightarrow G_N(p) = |g(p)|^N$

The cumulant generating function $\Omega_N(p) \equiv \ln G_N(p) = N \ln |g(p)|$
 $= \sum \frac{(-p)^k}{k!} \underbrace{\langle X^k \rangle}_{\text{cumulants of } S} = N w(p)$

If we introduce $w(p) = \sum \frac{(-p)^k}{k!} c_k \Rightarrow \boxed{C_k = N c_k}$ A8
 $\varphi_k^{1/k} = N^{1/k} c_k^{1/k} \sim N^{1/k}$

$\left. \begin{array}{l} \varphi_1 = N \langle x \rangle \rightarrow \text{trivial} \\ \varphi_2 = N \text{var}(x) \rightarrow \text{not so obvious} \end{array} \right\} \text{hyp 3}$

\rightarrow This implies $\boxed{\sigma_S = \sqrt{N} \sigma_x}$ \rightarrow typical mean deviation of S behaves as \sqrt{N} i.e. increases with N

Let us show that the higher cumulants are negligible from w_S \leftarrow All cumulants increase linearly with $N \Rightarrow C_k \propto N$

• We have seen above that the 2nd cumulant C_2 contributes to the fluctuation of S as $\sqrt{C_2} \rightarrow$

• By dimensional analysis, the $(k > 2)$ cumulant contributes to the fluctuation of S as $(C_k)^{1/k} \propto N^{1/k}$

$$\text{Therefore } \left(\frac{C_k^{1/k}}{C_2^{1/2}} \right) \propto N^{-\frac{(k-2)}{2k}} \rightarrow 0$$

$$\propto N^{\frac{1}{2} - \frac{1}{k}} = N^{-\frac{(k-2)}{2k}}$$

Therefore, in the large N limit

$$G_N(p) \simeq e^{-N c_1 p + \frac{1}{2} N c_2 p^2}$$

i.e.

$$\boxed{P_N(S) \underset{N \gg 1}{\simeq} \frac{1}{\sqrt{2\pi N \text{var} x}} \exp\left\{-\frac{(S - N \langle x \rangle)^2}{2N \text{var} x}\right\}}$$

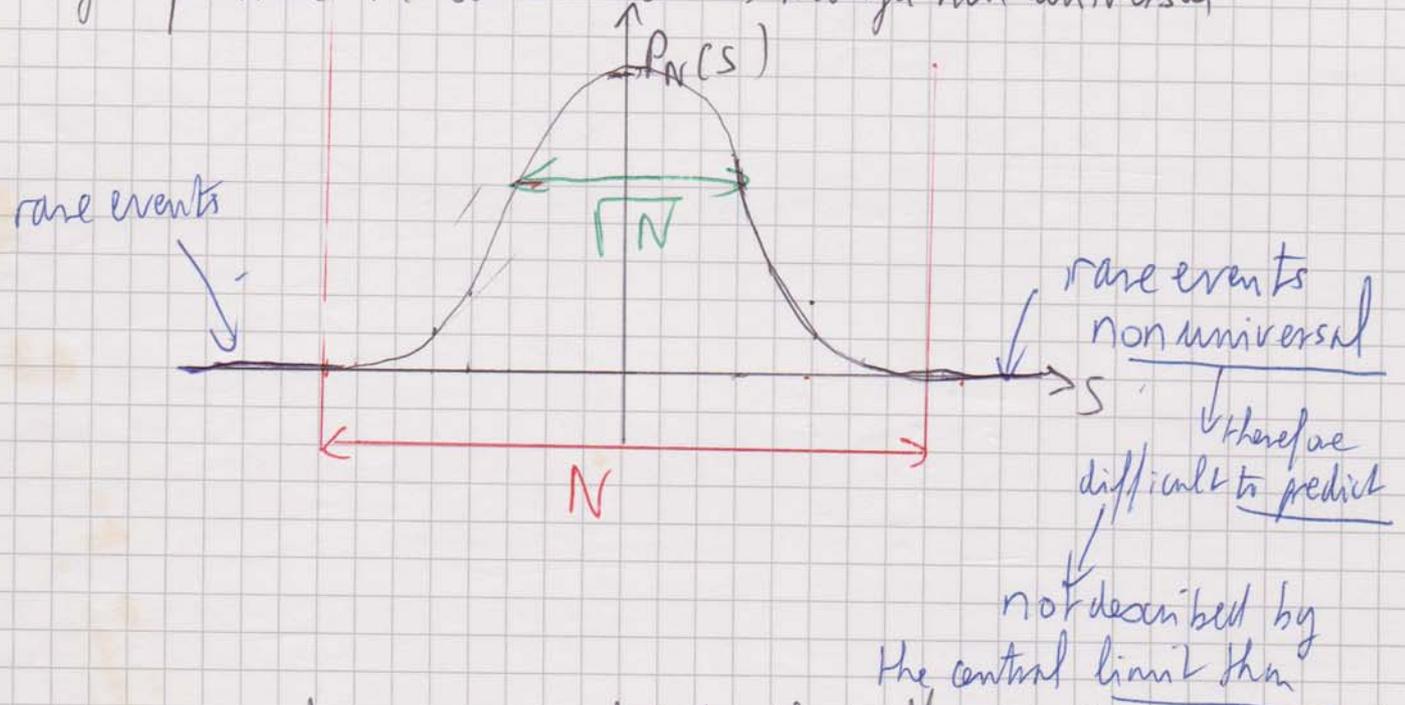
In the large N limit, S tends towards a Gaussian distribution function of mean $\langle S \rangle = N \langle x \rangle$ and standard deviation $\sigma_S = \sqrt{N} \sigma_x$

\Rightarrow This is remarkable because we only make very few hypotheses independent of $p(x)$!

\Rightarrow Emergence of a new universal law in the large N independent of microscopic details

This then is fundamental and actually perfectly illustrates that the laws governing the macroscopic (thermodynamic) may be completely \neq from the microscopic \rightarrow Depends on only few parameters (two here)

RQ: σ_s characterizes the standard deviations. This does not preclude the existence of rare events (or large deviations) which are also of importance in some cases \Rightarrow though non universal



\Rightarrow To analyze in more details the differences between universal fluctuations and rare events see Annex of Christophe book