

2.4) Ergodicity for a sphere in a fluid

$$f(u) = v(u)$$

$$[m\ddot{v}(t)] = -\gamma \dot{v}(t) - kx(t) + F(t) \quad \langle F(t) \rangle = 0$$

$$T = 2\pi \sqrt{\frac{m}{k}} \quad \tau_{\text{relax}} = m\gamma \ll T \rightarrow \text{strong friction}$$

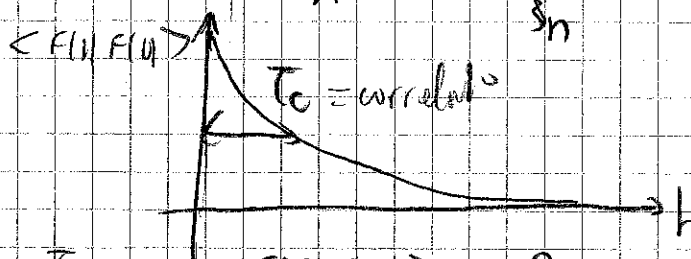
A) $\psi^2 = 0 \Rightarrow V_n = v(nz)$ with $z \ll \tau_{\text{relax}}$

$$\dot{v} = \frac{V_{n+1} - V_n}{T}$$

$$m \dot{v} = m \left(\frac{V_{n+1} - V_n}{\tau} \right) = -\gamma v_n + F(n\tau)$$

$$= V_{nt} = V_n \left(1 - \underbrace{\frac{\gamma \tau}{m}}_{\lambda} \right) + \underbrace{\frac{\tau}{m} F_n}_{\sum_n} = \lambda V_n + \sum_n$$

2) Typically



ii) $T \gg T_c$ $\langle F(T) A(u) \rangle \approx 0$, decoupled

$L, \tau_{\text{relax}} \gg \tau \gg \tau_c$ hierarchy of time scale

In this limit $\langle F(nz) | F(mz) \rangle \propto \delta_{nm}$

3) $\langle F^2(\tau) \rangle = \frac{\sigma^2}{\tau} \rightarrow$ Gaussian distrib^o

$$V_{n+1} = d V_n + \int_n = \dots = d^{n+1} V_0 + \sum_{k=0}^n d^k \int_{n-k} \rightarrow \text{Eq. (3)}$$

$$\text{var}(V_n) = \langle V_n^2 \rangle - \langle V_n \rangle^2 = d^{2n} \text{var}(V_0) + d^n \sum_{\ell} d^{\ell} \langle V_n s_{n,\ell} \rangle + \left(\sum_{\ell, \ell'} d^{\ell+\ell'} \langle s_{n,\ell} s_{n,\ell'} \rangle \right)$$

$$\text{var}(V_n) = \lambda^{2n} \text{var}(V_0) + \sum_{k=0}^{n-1} \lambda^{2k} \cdot \left(\frac{\tau}{m}\right)^2 \frac{\sigma^2}{\tau}$$

$$\text{var}(V_n) = \lambda^{2n} \text{var}(V_0) + \frac{1 - \lambda^{2n}}{1 - \lambda^2} \frac{\sigma^2 \tau}{m^2}$$

$$\lambda = 1 - \gamma \frac{\tau}{m} = 1 - \frac{\tau}{\tau_{\text{relax}}} < 1 \quad \left. \begin{array}{l} \lambda^{2n} \rightarrow 0 \\ \Rightarrow \text{var}(V_n) \rightarrow \frac{\sigma^2 \tau}{m^2} \frac{1}{1 - (1 - \frac{\tau}{\tau_{\text{relax}}})^2} \end{array} \right\} \Rightarrow \text{var}(V_n) \approx \frac{\sigma^2}{2m\gamma} \quad \left\| \begin{array}{l} \text{indep of} \\ \text{initial cond.} \end{array} \right.$$

$\Rightarrow V_n$ becomes stationary and also Gaussian

$$Pr(v) \xrightarrow{t \gg \tau_{\text{relax}}} P_{\text{eq}}(v) \text{ where } P_{\text{eq}}(v) = \frac{1}{\sqrt{2\pi} \sigma_v} e^{-\frac{v^2}{2\sigma_v^2}}$$

4) Fluctuation-dissipation

$$\left. \begin{array}{l} \langle E_c \rangle = \frac{1}{2} k_B T \\ \frac{1}{2} m \langle v^2 \rangle \end{array} \right\} \langle v^2 \rangle = \frac{k_B T}{m} = \frac{\sigma^2}{2m\gamma} \Rightarrow \boxed{\frac{\sigma^2}{2\gamma} = k_B T} \quad \begin{array}{l} \text{temperature} \\ \text{dissipated} \end{array}$$

5) We have $\sigma_v^2 = \frac{k_B T}{m} = \frac{\sigma^2}{2m\gamma} \Rightarrow P_{\text{eq}} \propto e^{-\frac{v^2}{2\sigma_v^2}} = e^{-\frac{E_c}{k_B T}} \quad \left\| \begin{array}{l} \text{Maxwell} \\ \text{distribution} \\ \text{func} \end{array} \right.$

Positron relaxation

overdamped $m \ddot{x} = -\gamma \dot{x} - kx + F(t) \Rightarrow \dot{x} = \frac{\gamma}{m} x + \frac{1}{m} F(t)$

we solve $\dot{v} = -\frac{\gamma}{m} v + \frac{1}{m} F(t)$ (same type of eq. here!)

Therefore substitute $\gamma \rightarrow k$
 $m \rightarrow \gamma$

$$\Rightarrow \text{var}(x(t)) = \frac{\sigma^2}{2\gamma k}$$

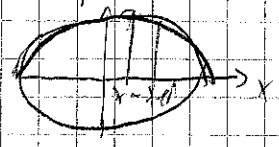
$$\sigma_x^2 = \text{var}(x(t)) = \frac{k_B T}{k} \Rightarrow P_{\text{eq}}(x) \propto e^{-\frac{E_p(x)}{k_B T}}$$

$$\left[P_{\text{eq}}(x|p) \propto P_{\text{eq}}(x) P_{\text{eq}}(p) \propto e^{-\frac{(E_c + E_p)/k_B T}{k_B T}} \right]$$

2.5) A) 1) $H(x, p) = \frac{p^2}{2m} + \frac{1}{2} m \omega^2 x^2$

$\dot{x} = \frac{\partial H}{\partial p} = \frac{p}{m} ; \dot{p} = -\frac{\partial H}{\partial x} = -m \omega^2 x$

2) $\ddot{x} = -m \omega^2 x \Rightarrow \ddot{x} + \omega^2 x = 0 \Rightarrow \dot{p} + \omega^2 p = 0 \Rightarrow \dot{x} + \omega^2 x = 0$



$E = \frac{1}{2} m \omega^2 x^2 + \frac{1}{2m} p^2$ - ellipse
 $\frac{1}{2} m \omega^2 x_0^2$ conserved
 $\Rightarrow x(t) = x_0 \cos(\omega t)$
 $p = -m x_0 \omega \sin(\omega t)$

3)

$\frac{dx}{dt} = -\omega x_0 \sin(\omega t) \Rightarrow dx = -\omega x_0 \sin(\omega t) dt = \frac{p}{m} dt$

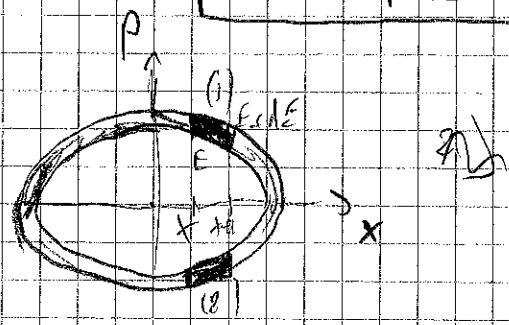
$\Rightarrow dt = \frac{m}{p} dx = \frac{m}{\sqrt{2mE - m^2 \omega^2 x^2}} dx$

!

It takes a time $T = \frac{dt}{\omega}$ for 1 cycle $\rightarrow$ Be careful, pass twice
 $\Rightarrow$ therefore $\frac{dt}{T} = \omega(1) dx \Rightarrow \omega(1) dx = \frac{m \omega}{\pi \sqrt{2mE - m^2 \omega^2 x^2}} dx$

$\Rightarrow \omega(x) = \frac{m \omega}{\pi \sqrt{2mE - m^2 \omega^2 x^2}}$

B) Stat Ψ
 1)



2)

$P_x dx = \frac{2}{\pi} \frac{E}{\omega}$

Area of a elliptical ring $A = \pi a b$ for an ellipse defined $\frac{x^2}{a^2} + \frac{y^2}{b^2} = 1$

we have $\frac{p^2}{2mE} + \frac{x^2}{2mE/m\omega^2} = 1$

$a = \sqrt{2mE} ; b = \sqrt{\frac{2E}{m\omega^2}}$

$A = \frac{2\pi E}{\omega} = E \times T \Rightarrow \boxed{\delta A = \delta E \times T}$

$$P(x)dx = \frac{2 \frac{dx}{\hbar} dp}{T dE} = \frac{2 dx dp}{T dE} = \frac{2 dx \left(\frac{m}{\hbar}\right) dE}{T dE}$$

④

$$P(x)dx = \frac{m \omega dx}{\pi \sqrt{2mE - m^2 \omega^2 x^2}} = \omega(x) dx$$

Same result!

Time average = Stat average over phase space!