

Exercise B

$$\begin{aligned}
 1) \quad \hat{H}_1 &= -\vec{\mu} \cdot \vec{B} = -\gamma B_x(t) \vec{S}_x - \gamma B_y(t) \vec{S}_y - \gamma B_z \vec{S}_z \\
 &= -\gamma B_1 \frac{\hbar}{2} (\cos(\omega t) \hat{\sigma}_x + \sin(\omega t) \hat{\sigma}_y) + \\
 &\quad \underbrace{-\gamma B_0 \frac{\hbar}{2} \hat{\sigma}_z}_{\omega_0}
 \end{aligned}$$

$\omega_1 \swarrow$
 $\omega_0 \swarrow$

Rappel: $\sigma_x = \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix}$ $\sigma_y = \begin{pmatrix} 0 & -i \\ i & 0 \end{pmatrix}$ $\sigma_z = \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix}$

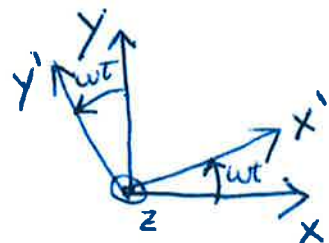
$$\Rightarrow \hat{H}_1 = \frac{\hbar}{2} \begin{pmatrix} \omega_0 & \omega_1 e^{-i\omega t} \\ \omega_1 e^{i\omega t} & -\omega_0 \end{pmatrix}; \quad \omega_0, \omega_1, \omega \geq 0$$

$$2) \quad V(t) = e^{+\frac{i}{\hbar} S_z \omega t} = \begin{pmatrix} e^{+i\omega t/2} & 0 \\ 0 & e^{-i\omega t/2} \end{pmatrix}$$

Example of the action of $V(t)$

$$\frac{\hbar}{2} \hat{\sigma}'_x = V(t) \frac{\hbar}{2} \hat{\sigma}_x V(t)^\dagger$$

$\hat{S}_x$ in the laboratory frame of reference / representation



The operator $\hat{S}_x$ described in the rotating frame of reference.

Be careful!

$$\frac{\hbar}{2} \hat{\sigma}'_x \neq \frac{\hbar}{2} \hat{\sigma}_x$$

In the rotating frame of reference
 $\hat{\sigma}'_x = \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix}$

Explicit calculation:

(3)

$$\hat{\sigma}_x' = \hat{V}(t) \hat{\sigma}_x \hat{V}(t)^\dagger = \begin{pmatrix} e^{i\omega t/2} & 0 \\ 0 & e^{-i\omega t/2} \end{pmatrix} \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix} \begin{pmatrix} e^{-i\omega t/2} & 0 \\ 0 & e^{i\omega t/2} \end{pmatrix} =$$

$$= \begin{pmatrix} 0 & e^{i\omega t} \\ e^{-i\omega t} & 0 \end{pmatrix}$$

Example $\omega t = \frac{\pi}{2} \rightarrow \hat{\sigma}_x' = -\hat{\sigma}_y$ CORRECT!

$\omega t = \pi \rightarrow \hat{\sigma}_x' = -\hat{\sigma}_x$ CORRECT!

$\omega t = \frac{3\pi}{2} \rightarrow \hat{\sigma}_x' = \hat{\sigma}_y$ CORRECT!

3) $\hat{H}_2 = \hat{V}(t) \hat{H}_1(t) \hat{V}(t)^\dagger + i\hbar \frac{d\hat{V}(t)}{dt} \hat{V}(t)^\dagger$

$\frac{d\hat{V}(t)}{dt} = \frac{i}{\hbar} \hat{S}^z \omega \hat{V}(t) \Rightarrow i\hbar \frac{d\hat{V}(t)}{dt} \hat{V}(t)^\dagger = -\omega \hat{S}_z = -\frac{\hbar\omega}{2} \hat{\sigma}_z$

$\hat{V}(t) \hat{H}_1(t) \hat{V}(t)^\dagger = \begin{pmatrix} e^{i\omega t/2} & 0 \\ 0 & e^{-i\omega t/2} \end{pmatrix} \frac{\hbar}{2} \begin{pmatrix} \omega_0 & \omega_1 e^{-i\omega t} \\ \omega_1 e^{i\omega t} & -\omega_0 \end{pmatrix} \begin{pmatrix} e^{-i\omega t/2} & 0 \\ 0 & e^{i\omega t/2} \end{pmatrix} =$

$$= \frac{\hbar}{2} \begin{pmatrix} \omega_0 e^{i\omega t/2} & \omega_1 e^{-i\omega t/2} \\ \omega_1 e^{i\omega t/2} & -\omega_0 e^{-i\omega t/2} \end{pmatrix} \begin{pmatrix} e^{-i\omega t/2} & 0 \\ 0 & e^{i\omega t/2} \end{pmatrix} =$$

$$= \frac{\hbar}{2} \begin{pmatrix} \omega_0 & \omega_1 \\ \omega_1 & -\omega_0 \end{pmatrix}$$

$\hat{H}_2 = \frac{\hbar}{2} \begin{pmatrix} \omega_0 - \omega & \omega_1 \\ \omega_1 & \omega - \omega_0 \end{pmatrix}$

$$4) \quad |\Psi_1(0)\rangle = \begin{pmatrix} 0 \\ 1 \end{pmatrix} \longrightarrow |\Psi_2(0)\rangle = \begin{pmatrix} 0 \\ 1 \end{pmatrix}$$

$\uparrow$ Laboratory Representation R_1
 $\uparrow$ Rotating frame of reference Representation R_2 .

$$|\Psi_2(t)\rangle = e^{-\frac{i}{\hbar} \hat{H}_2 t} |\Psi_2(0)\rangle$$

$$\text{S. } \omega = \omega_0: \quad H_2 = \frac{\hbar \omega_1}{2} \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix}$$

$$e^{-\frac{i}{\hbar} \hat{H}_2 t} = \sum_{K=0}^{+\infty} \frac{\left(-\frac{i}{\hbar} t \cdot \frac{\hbar \omega_1}{2}\right)^K}{K!} \left[\hat{\sigma}_x\right]^K \quad \text{Remark: } \sigma_x^2 = \mathbb{1}$$

$\hookrightarrow \hat{\sigma}_x$ in the representation R_2

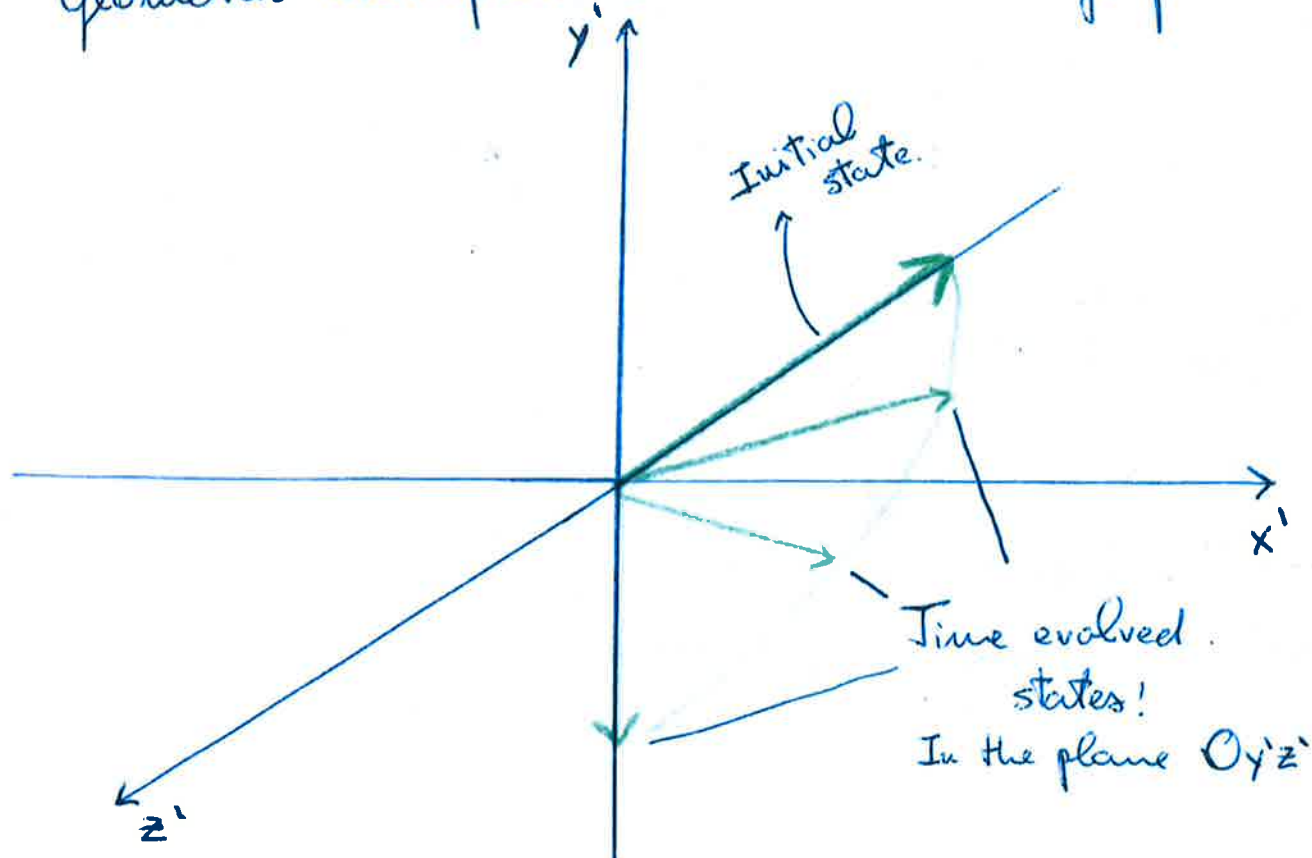
$$= \mathbb{1} \cdot \sum_{K \text{ even}} \frac{(-1)^{K/2} \left(\frac{\omega_1 t}{2}\right)^K}{K!} + i \hat{\sigma}_x \sum_{K \text{ odd}} \frac{(-1)^{\frac{K-1}{2}} \left(\frac{\omega_1 t}{2}\right)^K}{K!} =$$

$$= \mathbb{1} \cdot \cos\left(\frac{\omega_1 t}{2}\right) - i \hat{\sigma}_x \sin\left(\frac{\omega_1 t}{2}\right)$$

$$= \begin{pmatrix} \cos\left(\frac{\omega_1 t}{2}\right) & -i \sin\left(\frac{\omega_1 t}{2}\right) \\ -i \sin\left(\frac{\omega_1 t}{2}\right) & \cos\left(\frac{\omega_1 t}{2}\right) \end{pmatrix}$$

$$|\Psi_2(t)\rangle = \begin{pmatrix} -i \sin\left(\frac{\omega_1 t}{2}\right) \\ \cos\left(\frac{\omega_1 t}{2}\right) \end{pmatrix}$$

Geometric interpretation in the rotating frame. ⑤



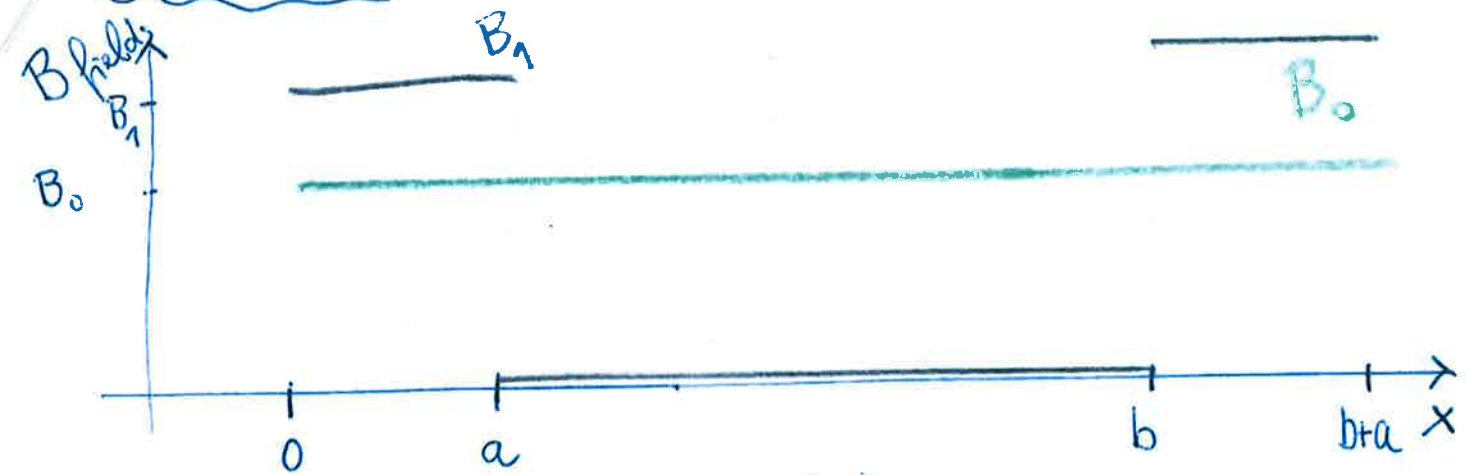
The time evolution is a rotation around the axis x' with angular velocity $-\omega_1$ (check for $\omega_1 t = \frac{\pi}{2}$, eigenstate of $\frac{\hbar\omega_1}{2}\sigma_y$, $-\frac{\hbar}{2}$)

Be CAREFUL: The axis x' is rotating with angular velocity ω around the axis $z \equiv z'$

$$\begin{aligned}
 5) \quad |\Psi_1(t)\rangle &= \hat{V}^\dagger(t) |\Psi_2(t)\rangle = \\
 &= \begin{pmatrix} e^{-i\omega_1 t/2} & 0 \\ 0 & e^{i\omega_1 t/2} \end{pmatrix} \begin{pmatrix} -i \sin \frac{\omega_1 t}{2} \\ \cos \frac{\omega_1 t}{2} \end{pmatrix} = \\
 &= \begin{pmatrix} -i e^{-i\omega_1 t/2} \sin \frac{\omega_1 t}{2} \\ e^{i\omega_1 t/2} \cos \frac{\omega_1 t}{2} \end{pmatrix}
 \end{aligned}$$

Exercice C

6



$$t=0, x=0 \quad |\Psi_1(t=0)\rangle = \begin{pmatrix} 0 \\ 1 \end{pmatrix} \quad \{\text{Laboratory Representation}\}$$

~~Wahlst-07~~

$$1) \quad t = \frac{a}{v} \Rightarrow x = a$$

We have to take into account the presence of the oscillating magnetic field. We use the results of exercise B and introduce the same rotating reference frame and the related representation R_2 .

$$|\Psi_2(t=0)\rangle = \begin{pmatrix} 0 \\ 1 \end{pmatrix} \quad \{\text{Rotating representation}\}$$

$$|\Psi_2(t=\frac{a}{v})\rangle = \begin{pmatrix} i \sin\left(\frac{\omega_1 a}{2v}\right) \\ \cos\left(\frac{\omega_1 a}{2v}\right) \end{pmatrix} \underset{\uparrow}{=} \frac{1}{\sqrt{2}} \begin{pmatrix} i \\ 1 \end{pmatrix}$$

Data of the exercise

$$\frac{\omega_1 a}{2v} = \frac{\pi}{4}$$

$$|\Psi_1(t=\frac{a}{v})\rangle = V(t)^\dagger |\Psi_2(t=\frac{a}{v})\rangle = \frac{1}{\sqrt{2}} \begin{pmatrix} i e^{-i\omega_1 a/2v} \\ e^{i\omega_1 a/2v} \end{pmatrix}$$

- 2) In both reference frames, the time evolution between a and b is simple.

Rotating reference: $H_2 = 0$ (we're talking $\omega = \omega_c$)

$$|\Psi_2(t = \frac{b}{v})\rangle = |\Psi_2(t = \frac{a}{v})\rangle = \frac{1}{\sqrt{2}} \begin{pmatrix} i \\ 1 \end{pmatrix}$$

Laboratory $H_1 = \frac{\hbar}{2} \begin{pmatrix} \omega_c & 0 \\ 0 & -\omega_c \end{pmatrix} = \frac{\hbar}{2} \begin{pmatrix} \omega & 0 \\ 0 & -\omega \end{pmatrix}$

$$\begin{aligned} |\Psi_1(t = \frac{b}{v})\rangle &= e^{-i \frac{1}{\hbar} \frac{\hbar}{2} \omega \frac{(b-a)}{v} \hat{\sigma}_z} |\Psi_1(t = \frac{a}{v})\rangle = \\ &= \frac{1}{\sqrt{2}} \begin{pmatrix} i e^{-i \omega b / (2v)} \\ e^{i \omega b / (2v)} \end{pmatrix} = V(t)^\dagger |\Psi_2(t = \frac{b}{v})\rangle \end{aligned}$$

- 3) Rotating reference: We use previous results.

$$\begin{aligned} |\Psi_2(t = \frac{b+a}{v})\rangle &= e^{-i \frac{1}{\hbar} \hat{H}_2 \frac{a}{v}} |\Psi_2(t = \frac{b}{v})\rangle = \\ &= \begin{pmatrix} \cos\left(\frac{\omega_1 a}{2v}\right) & i \sin\left(\frac{\omega_1 a}{2v}\right) \\ i \sin\left(\frac{\omega_1 a}{2v}\right) & \cos\left(\frac{\omega_1 a}{2v}\right) \end{pmatrix} \begin{pmatrix} i/\sqrt{2} \\ 1/\sqrt{2} \end{pmatrix} = \end{aligned}$$

$\frac{\omega_1 a}{2v} = \frac{\pi}{4}$ $\Downarrow$

$$\begin{pmatrix} 1/\sqrt{2} & i/\sqrt{2} \\ i/\sqrt{2} & 1/\sqrt{2} \end{pmatrix} \begin{pmatrix} i/\sqrt{2} \\ 1/\sqrt{2} \end{pmatrix} = \begin{pmatrix} i \\ 0 \end{pmatrix}$$

= eigenstate of $\hat{S}_z \equiv \hat{S}_{z'}$, $+\frac{\hbar}{2}$

Laboratory

(8)

$$|\Psi_2(t = \frac{b+a}{v})\rangle = V_2(t, \frac{b+a}{v})^\dagger |\Psi_1(t = \frac{b+a}{v})\rangle = i e^{-i\omega(b+a)/2v} \cdot \begin{pmatrix} 1 \\ 0 \end{pmatrix}$$

Probability of detecting $|\uparrow_z\rangle = 1$

4) $|\omega - \omega_0| \ll 1$ but $\omega \neq \omega_0$.

- The calculation of the ^{time-evolution} ~~passage~~ of the neutrons in the first region $x \in [0, a]$ goes as before because:

$$H_2 = \frac{\hbar}{2} \begin{pmatrix} \omega_0 - \omega & \omega_1 \\ \omega_1 & \omega - \omega_0 \end{pmatrix} \approx \frac{\hbar}{2} \begin{pmatrix} 0 & \omega_1 \\ \omega_1 & 0 \end{pmatrix}$$

Thus ~~$|\Psi_2(t = \frac{a}{v})\rangle$~~ $|\Psi_2(t = \frac{a}{v})\rangle = \frac{1}{\sqrt{2}} \begin{pmatrix} i \\ 1 \end{pmatrix}$

- The calculation of the ~~pass~~ time-evolution of the neutrons in the second region ~~the~~ $x \in [a, b]$ changes:

$$H_2 = \frac{\hbar}{2} \begin{pmatrix} \omega_0 - \omega & 0 \\ 0 & \omega - \omega_0 \end{pmatrix}$$

$$|\Psi_2(t = \frac{b}{v})\rangle = e^{-i \frac{\hbar}{2} H_2 \frac{b-a}{v}} \begin{pmatrix} i/\sqrt{2} \\ 1/\sqrt{2} \end{pmatrix} = \frac{1}{\sqrt{2}} \begin{pmatrix} i e^{-i(\omega_0 - \omega) \frac{b-a}{2v}} \\ e^{i(\omega_0 - \omega) \frac{b-a}{2v}} \end{pmatrix}$$

- The calculation of the time-evolution in the region $x \in [b, b+a]$ is similar to before:

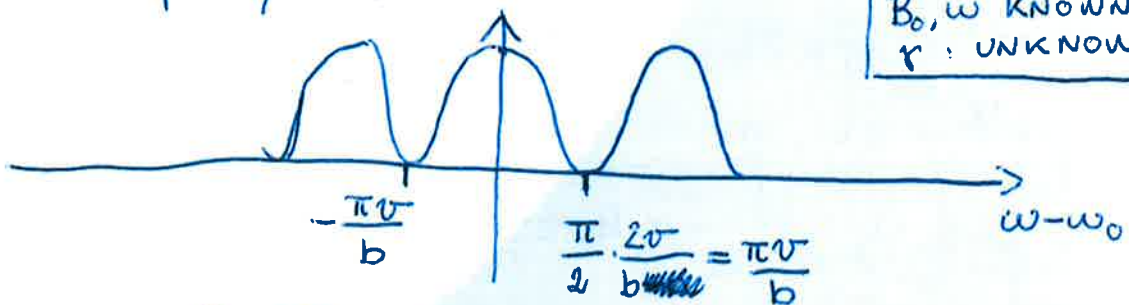
$$|\psi_2(t=\frac{b+a}{v})\rangle = \begin{pmatrix} 1/\sqrt{2} & i/\sqrt{2} \\ i/\sqrt{2} & 1/\sqrt{2} \end{pmatrix} \frac{1}{\sqrt{2}} \begin{pmatrix} i e^{-i(\omega_0-\omega)\frac{b-a}{2v}} \\ e^{i(\omega_0-\omega)\frac{b-a}{2v}} \end{pmatrix} =$$

$$= \frac{1}{2} \begin{pmatrix} i e^{-ix} + i e^{ix} \\ -e^{-ix} + e^{ix} \end{pmatrix} = \begin{pmatrix} i \cos x \\ i \sin x \end{pmatrix}$$

$$x = \frac{(\omega_0 - \omega)(b-a)}{2v}$$

$$\Rightarrow \text{Probability of measuring } |\uparrow_z\rangle = \cos^2\left(\frac{(\omega_0 - \omega)(b-a)}{2v}\right)$$

- Oscillatory dependence on $\omega - \omega_0$
- Recall that $\omega_0 = -\gamma B_0$, related to the intensity of B_z
 ω is the frequency of the rotation in the xy plane.
- For simplicity $b \gg a$



Largeur: $2 \times \frac{\pi v}{b}$

Typically:
 B_0, ω KNOWN
 γ : UNKNOWN



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We can interpret the width of the central peak as ~~an error~~ the fundamental uncertainty in the measurement of ω .

(10)

Quantum mechanics: A FREQUENCY (energy) measurement is more precise if it takes a lot of time, in this case the measurement time is b/ν . OK!

6) Δp : incertitude for the measurement of p .
 $\Delta \omega$: incertitude " " " " of ω .

$$\Delta p = \frac{\Delta \omega}{B_0} \approx \frac{2\pi\nu}{B_0 b}$$

$$b = \frac{2\pi\nu}{B_0 \Delta p}$$

Data: $\Delta p = 8,8 \cdot 10^{-7} \frac{q}{M_p}$

$$\nu = \frac{h}{\lambda_{dB} M_p} \rightarrow \lambda_{dB} = 3,1 \cdot 10^{-9} \text{ m}$$

$$B_0 = 1 \text{ T}$$

$$b = \frac{2\pi \cdot 6,62 \cdot 10^{-34}}{3,1 \cdot 10^{-9} \cdot 2,88 \cdot 10^{-7} \cdot 1,6 \cdot 10^{-19} \text{ s}} \cdot \frac{\text{m}^2 \text{ Kg}}{\text{m}} \cdot \frac{1}{\text{m}} \cdot \frac{\text{C} \cdot \text{s}}{\text{Kg}} \cdot \frac{1}{\text{C}} =$$

$$= \frac{2\pi \cdot 6,62}{3,1 \cdot 8,8 \cdot 1,62} \cdot 10 \cdot \text{m} = \text{AB} \text{ m } 4,7 \text{ m}$$

