

$$\text{I. 1) } [A+B, C] = (A+B)C - C(A+B) = AC + BC - CA - CB = \\ = [A, C] + [B, C]$$

$$\blacksquare [AB, C] = ABC - CAB = ABC - ACB + ACB - CAB = \\ = A[B, C] + [A, C]B$$

2) by contrapositive: suppose $[A, B] \neq 0$ and $\exists \{\varphi_n\}$ complete orthonormal set
 s.t. $A\varphi_n = a_n \varphi_n$ and $B\varphi_n = b_n \varphi_n$.

Then $\forall \psi \in \mathcal{H}$, $\psi = \sum_n c_n \varphi_n$ and:

$$[A, B]\psi = (AB - BA) \sum_n c_n \varphi_n = A \left(\sum_n c_n b_n \varphi_n \right) - \\ - B \left(\sum_n c_n a_n \varphi_n \right) = \\ = \sum_n c_n a_n b_n \varphi_n - \sum_n c_n b_n a_n \varphi_n = 0.$$

- As it's true $\forall \psi$, $[A, B] = 0$ (contradiction) □

3) using test function $\psi(x)$,

$$[X^m, P]\psi = \cancel{x^m} x^m (-i\hbar \frac{d\psi}{dx}) + i\hbar \frac{d(x^m \psi)}{dx} = \\ = -i\hbar x^m \frac{d\psi}{dx} + i\hbar \left[m x^{m-1} \psi + x^m \frac{d\psi}{dx} \right] = i\hbar m x^{m-1} \psi$$

$$\text{So } [X^m, P] = i\hbar m x^{m-1}$$

$$4) [A, B] = 0 \Rightarrow AB = BA \Rightarrow A^m B = B A^m$$

$$e^{-iA} B = \sum_n \frac{(-iA)^n}{n!} B = B \sum_n \frac{(-iA)^n}{n!} = B e^{-iA} \Rightarrow [e^{-iA}, B] = 0$$

- As $U(t) = e^{-\frac{i}{\hbar} H t}$, if $[H, B] = 0$ then we deduce $[U(t), B] = 0$; i.e. B action is unaffected by temporal evolution (more in the following exercises).

II.1

(2)

1) $\{A\}$ est ECOL iff there's a unique orthonormal basis such that $A = \sum_k \alpha_k |\varphi_k\rangle\langle\varphi_k|$. So A has to be nondegenerate

$$2) P_i' = |\langle\varphi_i|\psi'\rangle|^2 = |\langle\varphi_i|T^\dagger T|\psi\rangle|^2 = \langle\psi|T^\dagger T|\varphi_i\rangle\langle\varphi_i|T^\dagger T|\psi\rangle = \\ P_i = |\langle\varphi_i|\psi\rangle|^2 = \langle\psi|\varphi_i\rangle\langle\varphi_i|\psi\rangle = \langle P_i \rangle_\psi = \langle T P_i T \rangle_{T\psi}$$

where $P_i = |\varphi_i\rangle\langle\varphi_i|$ projector on i 's eigenspace of A

3) Probabilities have to stay the same after T , so either $T^{-1} = T$ unitary and linear (sends $(T\varphi, T\psi) = (\varphi, \psi)$), either $T^\dagger = T$ and antilinear (sends $(T\varphi, T\psi) = (\psi, \varphi)$). Notably, temporal inversion operator is the latter kind (ANTIUNITARY TRANSFORMATION).

4. the result above is further formalized by Wigner's theorem

$$4) A \xrightarrow{T} A' \quad \langle A' \rangle_{\psi'} = \langle A \rangle_\psi$$

$$\Rightarrow \langle\psi|T^\dagger A' T|\psi\rangle = \langle\psi|A|\psi\rangle \quad \forall |\psi\rangle \Rightarrow A = T^\dagger A' T \\ \Rightarrow \boxed{A' = T A T^\dagger}$$

$$5) [A, T] = AT - TA = \underset{\substack{\uparrow \\ A'=A}}{T A T^\dagger} T - T \underset{\substack{\uparrow \\ T^\dagger T=I}}{A} = TA - TA = 0$$

$$6) T^\dagger T = I \Rightarrow (I + \frac{i}{\hbar} \varepsilon G^\dagger + \dots) (I - \frac{i}{\hbar} \varepsilon G + \dots) = I \\ \Rightarrow I - \frac{i}{\hbar} \varepsilon G^\dagger + \frac{i}{\hbar} \varepsilon G = I \quad (\text{ignoring } o(\varepsilon) \text{ terms}) \\ \Rightarrow \boxed{G^\dagger = G}$$

$$7) A' = T A T^\dagger = (I - \frac{i}{\hbar} \varepsilon G + \dots) A (I + \frac{i}{\hbar} \varepsilon G + \dots) = \\ = A + \frac{i}{\hbar} \varepsilon (AG - GA) = A + \frac{i}{\hbar} \varepsilon [A, G] = A - \frac{i}{\hbar} \varepsilon [G, A]$$

II.2

(3)

$$1) T^{-1}(a)T(a)|x\rangle = |x\rangle \quad \forall |x\rangle \Rightarrow$$

$$\Rightarrow T^{-1}(a)|x+a\rangle = |x\rangle, \quad T(-a)|x+a\rangle = |x+a-a\rangle = |x\rangle \quad \forall x$$

$$\Rightarrow T^{-1}(a) = T(-a) \quad (\text{and } T^{-1} = T^\dagger \text{ from unitarity})$$

$$2) X'|x\rangle = T(a)X T^\dagger(a)|x\rangle = T(a)X|x-a\rangle = T(a)(x-a)|x-a\rangle = \\ = (x-a)|x\rangle = X|x\rangle - aI|x\rangle \quad \forall |x\rangle$$

$$\Rightarrow X' = X - aI$$

- ~~the~~ the analog of this phenomenon is the difference between translating the objects ($T(a)|x\rangle$) and translating the reference frame ($T(a)X T^\dagger(a)$). The minus sign shows for consistency between the two pictures.

$$3) \text{ We have } X' = X - \frac{i}{\hbar} \varepsilon [G, X] + \dots$$

$$X' = X - aI, \quad \text{in particular with } a = \varepsilon$$

$$\Rightarrow X - \frac{i}{\hbar} \varepsilon [G, X] = X - \varepsilon I \Rightarrow \frac{i}{\hbar} \varepsilon [G, X] = \varepsilon I \Rightarrow$$

$$\Rightarrow [X, G] = i\hbar \quad \boxed{G \equiv P}$$

$$4) \text{ part. libre} \Rightarrow V(x) = 0 \quad \forall x. \quad H(x) = -\frac{\hbar}{2m} \nabla^2 = \frac{P^2}{2m} = H(x+a)$$

$$\Rightarrow [T(a), H] = 0, \quad \text{and } [T(a), U(t)] = 0 \quad (\text{first exercise})$$

$$\frac{d}{dt} \langle \psi(t) | T(a) | \psi(t) \rangle = \frac{d}{dt} \langle \psi(0) | U^\dagger(t) T(a) U(t) | \psi(0) \rangle =$$

$$= \frac{d}{dt} \langle \psi(0) | U^\dagger(t) U(t) T(a) | \psi(0) \rangle = \frac{d}{dt} \langle T(a) \rangle_{\psi_0} = 0$$

$$5) \frac{d}{dt} \langle \psi(t) | T(a) | \psi(t) \rangle = \frac{d}{dt} \langle \psi_0 | (1 - \frac{i}{\hbar} \epsilon P) | \psi_0 \rangle = 0$$

$$\Rightarrow \frac{d}{dt} \langle P \rangle_{\psi_0} = 0, \text{ i.e. } P \text{ is a constant of motion}$$

↳ classical momentum

I (5) (facultatif):

- We define $g(x) = e^{Ax} e^{Bx}$. We try to obtain a differential equation

involving $g'(x)$:

$$\begin{aligned} \frac{dg(x)}{dx} &= A e^{Ax} e^{Bx} + e^{Ax} B e^{Bx} = \\ &= (A + e^{Ax} B e^{-Ax}) e^{Ax} e^{Bx} = (A + e^{Ax} B e^{-Ax}) g(x) \end{aligned}$$

- Define $f(x) = e^{Ax} B e^{-Ax}$. Taylor says: $f(x) = \sum_{m=0}^{\infty} \frac{x^m}{m!} f^{(m)}|_{x=0}$

$$f^{(1)}(x) = e^{Ax} A B e^{-Ax} - e^{Ax} B A e^{-Ax} = e^{Ax} [A, B] e^{-Ax}$$

$$f^{(2)}(x) = e^{Ax} A [A, B] e^{-Ax} - e^{Ax} [A, B] A e^{-Ax} = e^{Ax} [A, [A, B]] e^{-Ax} = 0$$

↑
Hypothesis

$$\Rightarrow f(x) = f(0) + x f^{(1)}|_{x=0} = B + x [A, B]$$

- We substitute and obtain: $g'(x) = (A + B + x [A, B]) g(x)$

that has, as a solution, $e^{xA+B+\frac{1}{2}x^2[A,B]}$

- the theorem is obtained with $x=1$.