

MQ PARTIEL OCT 2021

OP en MQ

$$1. P_u = |u\rangle\langle u| = \frac{1}{2} \begin{pmatrix} 1 \\ -i \end{pmatrix} \begin{pmatrix} 1 & i \end{pmatrix} = \frac{1}{2} \begin{pmatrix} 1 & i \\ -i & 1 \end{pmatrix}$$

$$2. P_u^\dagger = P_u \quad \text{OK}$$

$$3. \begin{vmatrix} \frac{1}{2} - \lambda & \frac{i}{2} \\ \frac{-i}{2} & \frac{1}{2} - \lambda \end{vmatrix} = \left(\frac{1}{2} - \lambda\right)^2 - \frac{1}{4} = -\lambda(1-\lambda) \quad \begin{matrix} \rightarrow \lambda = 0 \\ \rightarrow \lambda = 1 \end{matrix}$$

$$\lambda = 1 \rightarrow |u\rangle$$

$$\lambda = 0 \rightarrow |v\rangle \perp |u\rangle = \frac{\cancel{1-i} + i(1-i)}{\cancel{2}} = \frac{1-i+i-i^2}{2} = \frac{1-i+i+1}{2} = \frac{2-i+i}{2} = \frac{2}{2} = 1$$

normée.

La Mesure en MQ

1. a et b

$$2. P(a) = |\langle a|\psi\rangle|^2 = |\alpha|^2$$

$$P(b) = |\langle b,1|\psi\rangle|^2 + |\langle b,2|\psi\rangle|^2 = |\beta_1|^2 + |\beta_2|^2$$

3. si a $\rightarrow |a\rangle$

si b $\rightarrow \beta_1 |b,1\rangle + \beta_2 |b,2\rangle$ + renormalisation

$$\hookrightarrow \frac{\beta_1 |b,1\rangle + \beta_2 |b,2\rangle}{\sqrt{|\beta_1|^2 + |\beta_2|^2}}$$

$$4) \langle A \rangle = \langle \psi | A | \psi \rangle = |\alpha|^2 a + (|\beta_1|^2 + |\beta_2|^2) b$$

DOUBLE PITS QUANTUM

Double pits sym.

$$1. H = H_0 + W = \begin{pmatrix} E_0 - A & \\ & E_0 \end{pmatrix}$$

$$2. \begin{vmatrix} E_0 - E & -A \\ -A & E_0 - E \end{vmatrix} = (E_0 - E)^2 - A^2 \\ = (E_0 - A - E)(E_0 + A - E)$$

$$\Rightarrow E_{\pm} = E_0 \pm A$$

$$3. E_+ \rightarrow |+\rangle = \frac{|\psi_g\rangle - |\psi_d\rangle}{\sqrt{2}}; E_- \rightarrow |-\rangle = \frac{|\psi_g\rangle + |\psi_d\rangle}{\sqrt{2}}$$

$$4) |\psi(0)\rangle = |\psi_g\rangle = \frac{|+\rangle + |-\rangle}{\sqrt{2}}$$

$$\hookrightarrow |\psi(t)\rangle = \frac{e^{-\frac{iE_+ t}{\hbar}} |+\rangle + e^{-\frac{iE_- t}{\hbar}} |-\rangle}{\sqrt{2}} \\ = \frac{e^{-\frac{iE_0 t}{\hbar}} e^{-\frac{iAt}{\hbar}} |+\rangle + e^{-\frac{iE_0 t}{\hbar}} e^{+\frac{iAt}{\hbar}} |-\rangle}{\sqrt{2}}$$

$$5) P_d = |\langle \psi_d | \psi(t) \rangle|^2 \\ = \frac{1}{2} \left| e^{-\frac{iAt}{\hbar}} \langle \psi_d | + \rangle + e^{+\frac{iAt}{\hbar}} \langle \psi_d | - \rangle \right|^2$$

$$\text{avec } \langle \psi_d | + \rangle = -\frac{1}{\sqrt{2}} \text{ et } \langle \psi_d | - \rangle = \frac{1}{\sqrt{2}};$$

$$P_d = \frac{1}{4} \left| -e^{-\frac{iAt}{\hbar}} + e^{+\frac{iAt}{\hbar}} \right|^2$$

$$P_d = \sin^2\left(\frac{At}{\hbar}\right)$$

$$\text{m } \frac{At}{\hbar} = \frac{\pi}{2} (\pi) \text{ m a } P_d = 1$$

Deux puits asymétriques

$$6. \quad H' = H_0' + W = \begin{pmatrix} E_g - A & \\ & -A \quad E_d \end{pmatrix}$$

$$7. \quad \begin{vmatrix} E_g - E & -A \\ -A & E_d - E \end{vmatrix} = (E_g - E)(E_d - E) - A^2$$

$$= E_g E_d - (E_g + E_d)E + E^2 - A^2$$

$$= E^2 - (E_g + E_d)E + E_g E_d - A^2$$

$$\hookrightarrow E_{\pm}' = \frac{1}{2} (E_g + E_d \pm \sqrt{(E_g - E_d)^2 + 4A^2})$$

$$8. \quad E_{\pm}' = \frac{1}{2} (E_g + E_d \pm |E_g - E_d| \sqrt{1 + \frac{4A^2}{(E_g - E_d)^2}})$$

$$\approx \frac{1}{2} (E_g + E_d \pm |E_g - E_d| + \frac{2A^2}{(E_g - E_d)})$$

$$E_{\pm}' \text{ en } A^2 \quad E_{\pm} \text{ en } A!$$

m. A petit, puits asym - sensible au couplage tunnel que le puits sym.

$$9) \quad |\psi(t=0)\rangle = |\psi_g\rangle = \cos(\frac{\theta}{2}) |+\rangle - \sin(\frac{\theta}{2}) |-\rangle$$

$$\hookrightarrow |\psi(t)\rangle = \cos(\frac{\theta}{2}) e^{-\frac{iE_+ t}{\hbar}} |+\rangle - \sin(\frac{\theta}{2}) e^{-\frac{iE_- t}{\hbar}} |-\rangle$$

$$10) \quad P_d = |\langle \psi_d | \psi(t) \rangle|^2$$

$$= \left| \cos(\frac{\theta}{2}) e^{-\frac{iE_+ t}{\hbar}} \langle \psi_d | + \rangle - \sin(\frac{\theta}{2}) e^{-\frac{iE_- t}{\hbar}} \langle \psi_d | - \rangle \right|^2$$

$$\text{avec } \langle \psi_d | + \rangle = \sin \frac{\theta}{2} \text{ et } \langle \psi_d | - \rangle = \cos \frac{\theta}{2}$$

$$\Rightarrow P_d = \left| \sin \frac{\theta}{2} \cos \frac{\theta}{2} \right|^2 \left| e^{-\frac{iE_+ t}{\hbar}} - e^{-\frac{iE_- t}{\hbar}} \right|^2$$

$$= \frac{1}{4} \sin^2 \theta \times 4 \sin^2 \left(\frac{(E_+ - E_-) t}{2\hbar} \right)$$

$$\hookrightarrow = \sin^2 \theta \sin^2 \left(\frac{(E_+ - E_-) t}{2\hbar} \right)$$

$$\sin^2 \theta = \frac{E_{in}^2 \theta}{1 + E_{in}^2 \theta} = \frac{4A^2 / \Delta E^2}{1 + \frac{4A^2}{\Delta E^2}} = \frac{4A^2}{4A^2 + \Delta E^2}$$

$$\text{et } E_+ - E_- = \sqrt{\Delta E^2 + 4A^2}$$

$$\Rightarrow P_{ul}(t) = \frac{4A^2}{4A^2 + \Delta E^2} \sin^2 \frac{\sqrt{\Delta E^2 + 4A^2} E}{2\hbar}$$

$$\text{en } g \text{ est } P_{0 \text{ MAX}} = \frac{4A^2}{4A^2 + \Delta E^2} < 1 \quad \text{donc jamais } 100\%$$

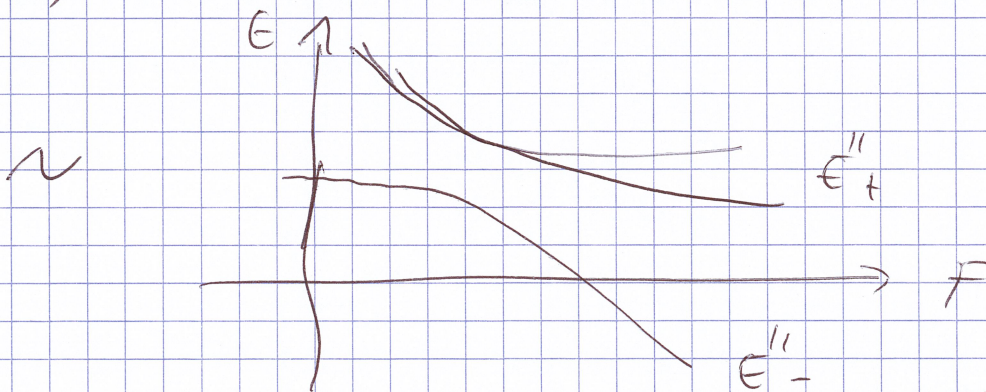
soit de trouver l'é⁻ en l'état (4d)

$$11) \quad H'' = \begin{pmatrix} E_g + \alpha_g F & -A \\ -A & E_d + \alpha_d F \end{pmatrix}$$

change^t $E_g \rightarrow E_g + \alpha_g F$ et $E_d \rightarrow E_d + \alpha_d F$

$$\text{il en } E_{\pm}'' = \frac{1}{2} \left(E_g + \alpha_g F + E_d + \alpha_d F \pm \sqrt{(E_g + \alpha_g F - E_d - \alpha_d F)^2 + 4A^2} \right)$$

12) 2 branches d'hyperboles



$$13) \quad S = E_+'' - E_-'' = \sqrt{(E_g - E_d + (\alpha_g - \alpha_d)F)^2 + 4A^2}$$

$$\sqrt{x^2 + a^2} \text{ min si } x = 0$$

$$\hookrightarrow S \text{ min si } (E_g - E_d + (\alpha_g - \alpha_d)F)^2 = 0$$

i.e. si $E_g + \alpha_g F = E_d + \alpha_d F$

~~14)~~ S et leurs $S_{\text{max}} = 2A$

14) "Fréq de Bohr" (ou plutôt pulsation)

$$\omega_{\text{Bohr}} = \frac{E_+ - E_-}{\hbar} = \frac{S}{\hbar} = \frac{2A}{\hbar} \text{ à résonance}$$

15) Dipôle oscillant \rightarrow rayonnement

$$\lambda = \frac{c}{\nu} = \frac{2\pi c}{\omega} = \frac{2\pi \hbar c}{2A}$$

$$\Rightarrow A = \pi \frac{\hbar c}{\lambda}$$

$$\text{AN } A \approx 3 \times \frac{200 \text{ MeV fm}}{10^{-10} \text{ m}} = 3 \times \frac{200 \cdot 10^6 \text{ eV fm}}{10^{-10} \text{ fm}}$$

$$A \approx 6 \times 10^3 \text{ eV}$$

Question subsidiaire:

énergies liaisons typiques des conducteurs

ou \sim conducteurs \sim eV

\hookrightarrow couplage faible