STUDENT SELF-ASSESSMENT Nonlinear wave equation & Nonlinear susceptibility tensors

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The electric field amplitude of a wave at ω_j , which propagates along the direction (Oz), is described by the relation : $\mathcal{E}_j(z,t) = \mathbf{E}(\omega_j)e^{-i\omega_j t} + C.C.$, with $\mathbf{E}(\omega_j) = A_j(z)e^{ik(\omega_j)z}\mathbf{e}_j$. The related field intensity is given by: $I_j = 2n(\omega_j)c \epsilon_0 |A_j(z)|^2$.

EXERCISE 1

A monochromatic electromagnetic wave is propagating in a nonlinear and homogenous medium. Its complex amplitude $E(\omega)$ satisfies the wave equation:

$$\boldsymbol{\nabla} \times \boldsymbol{\nabla} \times \boldsymbol{E}(\omega) = \frac{\omega^2}{c^2} \underline{\underline{\epsilon}}(\omega) \boldsymbol{E}(\omega) + \omega^2 \mu_0 \boldsymbol{P}_{NL}(\omega).$$
(1)

One reminds the relation: $\nabla \times \nabla \times E = \nabla(\nabla \cdot E) - \Delta E$. In the following, the specific case of an **isotropic medium** will be considered.

- 1. What can you say about $\underline{\epsilon}(\omega)$ and the term $\nabla \cdot E$? Justify your answer.
- 2. Assuming a wave propagating along the z direction, rewrite Eq. (1).
- 3. Expressing the complex amplitude as $\boldsymbol{E} = A(z)e^{ikz}\boldsymbol{e}$, show that the wave equation is reduced to a first order differential equation in z. Explicit any approximation(s).

EXERCISE 2

You are in charge of a new project in your company: the development of a laser source at 532 nm based on the PicoYAG laser module, already available in your product catalogue. The PicoYAG delivers sub-ns pulse duration at 1064 nm with a 1 kW peak-power.

You seek to evaluate an order of magnitude for the second-harmonic generation (SHG) efficiency throughout a 2 mm long lithium niobate (LiNbO₃) crystal. The refractive indices at 1064 nm and 532 nm are respectively equal to $n_{\omega} = 2.155$ and $n_{2\omega} = 2.234$. In the following, the effective nonlinear susceptibility is taken equal to 18 pm/V. The beam diameter inside the crystal is kept equal to 170 μ m for both beams at ω and 2 ω . Subsequently, the undepleted pump approximation is assumed.

- 1. Derive the phase matching condition. Can it be fulfilled?
- 2. Show that the intensity at 2ω follows

$$I_{2\omega,z} = \frac{(2\omega)^2}{2n^3 c^3 \varepsilon_0} \left| \chi_{eff}^{(2)} \right|^2 \sin^2 \left(\frac{\Delta k}{2} z \right) \frac{I_{\omega}^2}{\Delta k^2} \tag{2}$$

Comment on the evolution of $I_{2\omega,z}$ for various Δk situations.

- 3. For $\Delta k \neq 0$, show that the maximum SHG efficiency is reached for $L_c = \pi/\Delta k$, the coherence length. Calculate L_c and the maximum expected SHG efficiency.
- 4. The SHG efficiency relation (2) can be rewritten :

$$\eta_{\rm SHG} = \frac{I_{2\omega}}{I_{\omega}} = \frac{(2\omega)^2}{8\epsilon_0 n_{\omega}^2 n_{2\omega} c^3} \left| \chi_{\rm eff}^{(2)} \right|^2 \, {\rm sinc}^2 (\Delta kz/2) \, I_{\omega} \, z^2.$$

Assuming a perfect phase matching condition, calculate the expected SHG efficiency. Comment about the validity of the undepleted pump approximation.

 $\varepsilon_0 = \! 8.85 \times 10^{-12} \ \mathrm{F \ m^{-1}}$