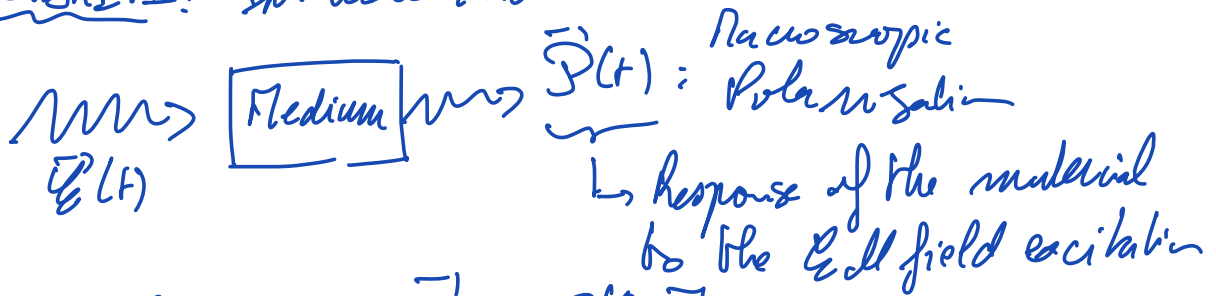


NONLINEAR OPTICS COURSE

LECTURE 1: INTRODUCTION



LINEAR REGIME: $\vec{P}(t) = \epsilon^{(1)} \vec{E}(t)$

Assuming a dielectric material.

NONLINEAR REGIME $\rightarrow$ The response of the material is no longer linear

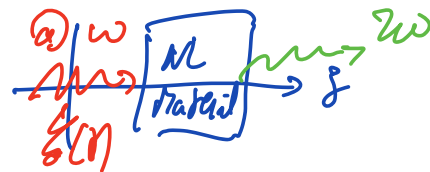
$$\vec{P}(t) = \epsilon^{(1)} \vec{E}(t) + \epsilon^{(2)} \vec{E}(t) \vec{E}(t) + \epsilon^{(3)} \vec{E}(t) \vec{E}(t) \vec{E}(t) + \dots$$

$\underbrace{\epsilon^{(2)} \vec{E}(t) \vec{E}(t)}_{\text{2nd order NL response}} + \underbrace{\epsilon^{(3)} \vec{E}(t) \vec{E}(t) \vec{E}(t)}_{\text{3rd order NL response}} + \dots$

ILLUSTRATION (1)

Monochromatic optical field incident on a 2nd-order NL material

$$\vec{E}(t) = A \cos(\omega t - k z)$$



$\Rightarrow$ 2nd-order NL response

$$\vec{P}^{(2)}(t) = \epsilon^{(2)} A^2 \cos^2(\omega t - k z)$$

$$\boxed{\vec{P}^{(2)}(t) = \frac{\epsilon^{(2)}}{2} A^2 \left[1 + \cos(2\omega t - 2kz) \right]}$$

$\underbrace{\frac{\epsilon^{(2)}}{2} A^2}_{\text{NL response } (\omega = 0)} + \underbrace{\cos(2\omega t - 2kz)}_{\omega = 2\omega}$

OPTICAL
RECTIFICATION

$\hat{L} \rightarrow$ 2nd order Harmonic
Generation
effect

ILLUSTRATION 2

$$\begin{matrix} \omega_1 \\ \omega_2 \end{matrix} \rightarrow \boxed{\text{2nd order NLI}} \rightarrow \mathcal{P}^{(2)}(t) = \mathcal{R} \{ E E \}$$

$$E(t) = A_1 \cos(\omega_1 t + \phi_1) + A_2 \cos(\omega_2 t + \phi_2)$$

$$R \{ E(t) = \underbrace{E(\omega_1)} e^{-i\omega_1 t} + \underbrace{E(\omega_2)} e^{-i\omega_2 t} + c.c.$$

$$\mathcal{P}^{(2)}(t) = \mathcal{R} \{ E^2(t) \} \quad \text{complex amplitudes}$$

$$R \{ \mathcal{P}^{(2)}(t) = P(0) e^{i(0)t} + P(2\omega_1) e^{-2i\omega_1 t} + P(2\omega_2) e^{-2i\omega_2 t} + \\ P(\omega_1 + \omega_2) e^{-i(\omega_1 + \omega_2)t} + P(\omega_1 - \omega_2) e^{-i(\omega_1 - \omega_2)t} + c.c.$$

with

$$\left\{ \begin{aligned} P(0) &= 2\mathcal{R}^{(2)} \{ E(\omega_1) E(-\omega_1) + E(\omega_2) E(-\omega_2) \} \\ &\quad \text{with } E(-\omega) = E^*(\omega) = \text{complex conjugate} \\ P(2\omega_{\frac{1}{2}}) &= \mathcal{R}^{(2)} \{ E(\omega_{\frac{1}{2}}) E(\omega_{\frac{1}{2}}) \} \\ P(\omega_1 \pm \omega_2) &= 2\mathcal{R}^{(2)} \{ E(\omega_1) E(\pm\omega_2) \} \end{aligned} \right.$$

CONCLUSIONS

$\hookrightarrow$ we understand that the interaction between 2 waves @ ω_1 and ω_2 through a 2nd order N

material can lead to the generation of waves:

@ $\omega = 0 \Rightarrow$ generation of a static field in the material.

@ $2\omega_1$ and $2\omega_2 \Rightarrow$ SHG

@ $\omega_3 = \omega_1 + \omega_2 \Rightarrow$ SUM FREQ. Generation

@ $\omega_3 = \frac{\omega_1 - \omega_2}{\omega_2 - \omega_1} \Rightarrow$ DIFFERENCE Freq. Generatn.

$\hookrightarrow$ The magnitude of the NL response is proportional to

• $\chi^{(2)}$: NL response of the material

• $E(\omega_1)E(\omega_2)$

$\Rightarrow$ NONLINEAR EFFECTS are governed by the magnitudes of the ELECTRIC FIELDS

Intensity $\propto |E|^2$

NL OPTICS is driven by FIELD INTENSITY